

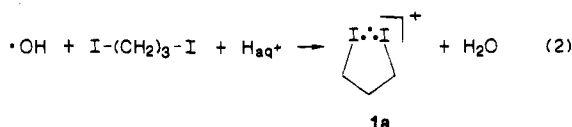
Table I. Properties of Intramolecular Radical Cations of Type 1<sup>a</sup>

$\left( \begin{array}{c} \cdot \text{I} \cdots \text{I} \cdot \\ \text{(CH}_2\text{)}_n \\ \mathbf{1} \end{array} \right)^+$	$\lambda_{\text{max}}$ , nm	$t_{1/2}$ , $\mu\text{s}$	$G$ , per 100 eV	$\epsilon$ , $\text{M}^{-1} \text{cm}^{-1}$
$n = 6$	455	75	0.6	5500
$n = 5$	445	85	1.7	4500
$n = 4$	435	120	2.1	4700
$n = 3$ ( <b>1a</b> )	405	150	2.9	4400
$n = 2$	(550)		(<0.3)	

<sup>a</sup> Experimental conditions: pulse radiolysis of N<sub>2</sub>O saturated, pH 3, aqueous solutions of  $2 \times 10^{-5}$  M 1,*n*-diiodoalkanes. Error limits:  $\pm 5$  nm for  $\lambda_{\text{max}}$ ,  $\pm 50$  nm for  $n = 2$ ;  $\pm 10\%$  for  $t_{1/2}$ ,  $G$ , and  $\epsilon$ ,  $\pm 20\%$  for  $n = 6$ .

•OH-induced oxidation of monoiodoalkanes, i.e., the radical cations are only formed at low pH (<5). The extinction coefficient of the transient is calculated to  $4400 \text{ M}^{-1} \text{cm}^{-1}$ .

The absorption is assigned to the intramolecular radical cation **1a** formed in the overall process



The alternative possibilities, namely a molecular radical cation without iodine-iodine interaction,  $\text{I}(\text{CH}_2)_3-\text{I}^{\cdot+}$ , or a bimolecular complex of type **2**, can be excluded. The former is unlikely because all other  $\text{RI}^{\cdot+}$  in aqueous solution have been found to absorb around 315 nm.<sup>5</sup> The dimer species, on the other hand, showed a pronounced dependence on the solute concentration with hardly any detectable yield ( $G < 0.2$ ) at the concentration of  $2 \times 10^{-5}$  M used in our present experiment.<sup>5</sup> Low solubilities ( $\leq 5 \times 10^{-5}$  M) unfortunately precluded the study of such a concentration dependence for the diiodo alkanes. A most important argument in support of **1a** is its lifetime which exceeds that of  $\text{RI}^{\cdot+}$  and  $(\text{RI} \cdots \text{IR})^{\cdot+}$  by almost one order of magnitude. The high stability of **1a** can, of course, be rationalized by its sterically most favorable five-membered ring structure.

Compatible results have been obtained in the oxidation of a number of other 1,*n*-diiodoalkanes. The measured data are summarized in Table I. It can be seen that an increasing chain length between the two iodine atoms and thus an increasing ring size of **1** is accompanied by a red-shift in  $\lambda_{\text{max}}$  and a decrease in lifetime and yield. All this matches with corresponding observations for numerous intramolecular S:S-bonded radical cations generated in the oxidation of organic dithia compounds.<sup>1</sup> The underlying rationale is a lower degree of hetero atom p-orbital overlap in the larger ring system.

A geometry situation which is even less favorable would exist in strained three- or four-membered ring structures as has been demonstrated in the oxidation of various sulfur compounds.<sup>1,9</sup> In our present case the oxidation of 1,2-diiodoethane would represent such an example. In fact, only a very small yield of radical cations seems to be formed ( $G < 0.3$  according to conductivity measurements) besides a larger yield of neutral, as yet unidentified radicals. The observable UV-absorption in a pulsed aqueous 1,2-diiodoethane solution shows, however, a long tailing with a shoulder around  $500 \pm 50$  nm. Assignment of such a red-shifted band to an intramolecular radical cation would be reasonable since the optical transition energies for  $2\sigma/1\sigma^*$  species are a direct function of orbital overlap.<sup>1,2</sup> It needs to be mentioned that a relatively fast hydrolysis of the 1,2-diiodoethane leaves, however, some ambiguity about this last result.

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## The Photosensitized Production of Singlet Molecular Oxygen ( $^1\Delta_g\text{O}_2$ ) in a Solid Organic Polymer Glass: A Direct Time-Resolved Study

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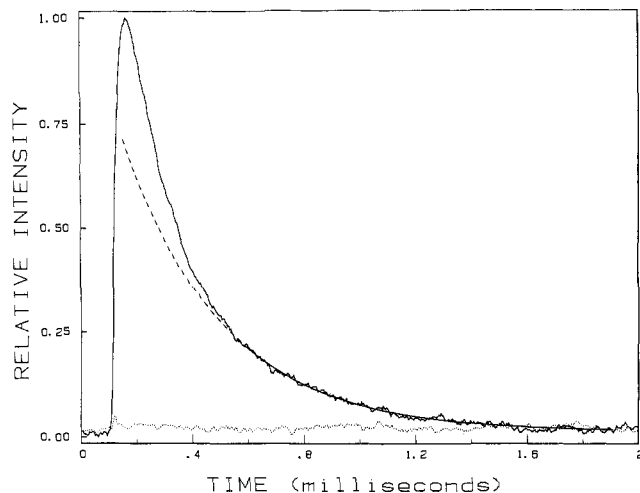
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The photosensitized production of singlet molecular oxygen ( $^1\Delta_g\text{O}_2$ ) has long been a subject of scientific study.<sup>1-6</sup> Recently, direct evidence for the presence of  $^1\Delta_g\text{O}_2$  in thin (10–30  $\mu\text{m}$ ) polymer films was obtained by Byteva et al.<sup>7,8</sup> who detected  $^1\Delta_g\text{O}_2$  phosphorescence ( $^3\Sigma_g^-\text{O}_2 \leftarrow ^1\Delta_g\text{O}_2$ ; 1270 nm) in a steady-state photosensitized experiment. In a time-resolved study,  $^1\Delta_g\text{O}_2$  has also been observed in a heterogeneous ionomer system (Nafion powders) in which the  $^1\Delta_g\text{O}_2$  sensitizer was apparently located at the surface of solvent filled or desiccated cavities.<sup>9</sup> We now report that  $^1\Delta_g\text{O}_2$  phosphorescence can be detected in a time-resolved photosensitized experiment from a homogeneous, solid polymethylmethacrylate (PMMA) glass. As in the solution phase,<sup>10-26</sup> we expect this technique will become an indispensable

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**Figure 1.** The time-resolved, near IR luminescence signal observed subsequent to irradiation of acridine ( $9.3 \times 10^{-4}$  M, 355 nm, 1.0 mJ pulse $^{-1}$ ) in an aerated, 22 °C, PMMA glass (solid line). The data shown are an average of 16 independent signals. Under these conditions, the half-life for  $^1\Delta_g\text{O}_2$  decay is approximately 240  $\mu\text{s}$ . A single exponential function has been applied to the latter portion of the signal. A substantial deviation from first-order decay kinetics is apparent. Near IR luminescence was not observed from a similar polymer sample prepared in the absence of oxygen (dotted line). The latter data are an average of 128 independent signals.

tool in studies of  $^1\Delta_g\text{O}_2$  photophysics and photochemistry. In particular, the role of  $^1\Delta_g\text{O}_2$  in the photooxidation and photodegradation of solid polymers has been a focus of debate.<sup>4-6</sup> Since we directly monitor molecular oxygen, this experimental approach also provides a unique method to accurately quantify oxygen diffusion in a polymer.<sup>27-30</sup>

PMMA samples<sup>31</sup> were prepared containing one of five different sensitizers.  $^1\Delta_g\text{O}_2$  was identified as the carrier of the observed signal<sup>32</sup> (Figure 1) on the basis of the following points: (1)  $\lambda_{\text{max}}$  of luminescence was  $\sim 1270$  nm, in agreement with previous  $^1\Delta_g\text{O}_2$  assignments.<sup>7,8,15,16,34</sup> (2) A sample prepared in the absence of oxygen failed to give a signal upon irradiation (Figure 1).<sup>35</sup> (3) Irradiation of a sensitizer free, oxygenated PMMA sample failed to yield a signal. (4) Samples prepared with the known  $^1\Delta_g\text{O}_2$  quencher, 1,4-diazabicyclo[2.2.2]octane (DABCO), led to an increased rate of signal decay.<sup>36</sup> (5) Polymer samples stored under air gave rise to a signal whose peak intensity was less than that observed for identical samples stored under atmospheres with a larger partial pressure of oxygen. More importantly, the rate of both signal appearance and disappearance increased as the oxygen concentration in the sample increased. We suggest that these observations reflect, respectively, an increased rate of sensitizer quenching by oxygen<sup>19</sup> and the contribution of an oxygen-induced  $^1\Delta_g\text{O}_2$  deactivation channel.<sup>36</sup>

The time-resolved  $^1\Delta_g\text{O}_2$  signals were characterized by three general features: (1) A substantial deviation from first-order decay

kinetics was observed in each case, (2) the rate of  $^1\Delta_g\text{O}_2$  decay was dependent both on the sensitizer and the sensitizer concentration, and (3) The half-lives ( $\tau_{1/2}$ ) for  $^1\Delta_g\text{O}_2$  decay and rise in air-saturated samples ( $\sim 300$  and  $\sim 10$   $\mu\text{s}$ , respectively) were approximately an order of magnitude larger than the corresponding values in an appropriate liquid phase PMMA monomer model, methyl propionate. As in the case of other liquid phase organic solvents,<sup>10-19</sup> the  $^1\Delta_g\text{O}_2$  decay in methyl propionate followed first-order kinetics ( $\tau_{\text{decay}} = 38$   $\mu\text{s}$ ).<sup>26,32,36</sup> In an attempt to quantify the PMMA data, we have applied a single exponential function to portions of the decay signal (Figure 1). This approach indicates clearly the extent to which  $^1\Delta_g\text{O}_2$  deactivation is determined by mixed-order kinetics. For the five sensitizers examined, the rate of  $^1\Delta_g\text{O}_2$  decay decreased in the order phenazine ( $\tau_{1/2} \sim 110$   $\mu\text{s}$ ) > acridine ( $\tau_{1/2} \sim 225$   $\mu\text{s}$ ) > fluoranthene and 1,12-benzoperylene ( $\tau_{1/2} \sim 260$   $\mu\text{s}$ ) > 2'-acetonaphthone ( $\tau_{1/2} \sim 425$   $\mu\text{s}$ ).<sup>37</sup> These rates reflect the relative efficiency with which the sensitizers quench  $^1\Delta_g\text{O}_2$  in solution.<sup>36</sup>

At this juncture, our results may be interpreted to support several different models. The kinetics of  $^1\Delta_g\text{O}_2$  appearance and disappearance, and the rate of signal change subsequent to a perturbation in the ambient atmosphere,<sup>35</sup> indicate that physical displacement of molecular oxygen in the polymer host, in the period defined by the  $^1\Delta_g\text{O}_2$  lifetime ( $\sim 100$   $\mu\text{s}$ ), is substantially less than the displacement allowed by diffusion in the solution phase.<sup>35,38</sup> As a result,  $^1\Delta_g\text{O}_2$  may encounter several different environments within its lifetime, each characterized by a unique set of  $^1\Delta_g\text{O}_2$  quenching conditions. For example, upon formation in a photosensitized process,  $^1\Delta_g\text{O}_2$  will necessarily be adjacent to the sensitizer which, in turn, may act as a quencher.<sup>15</sup> Diffusion of  $^1\Delta_g\text{O}_2$  will change both the oxygen-sensitizer distance and relative geometric orientation and may expose  $^1\Delta_g\text{O}_2$  to a slightly different bulk polymer environment.<sup>27-28</sup>  $^1\Delta_g\text{O}_2$  may also encounter a photoinduced transient (e.g., a sensitizer triplet state<sup>39</sup> or a polymer-derived intermediate) which, in turn, could give rise to a time dependent quenching channel.<sup>26,40</sup> Within the period of  $\sim 100$ – $200$   $\mu\text{s}$ , therefore, the relative probability of  $^1\Delta_g\text{O}_2$  quenching will change, and a deviation from first-order decay kinetics will result.<sup>41</sup> Alternatively, since a biexponential function can be used to fit the  $^1\Delta_g\text{O}_2$  decay, we must also consider a model for which evidence has been presented in the literature.<sup>46-48</sup> Small gas molecules dissolved in the polymer may be sorbed in two distinct modes, ordinary dissolution and adsorption in microvoids. Experiments currently in progress in our laboratory are designed to further characterize  $^1\Delta_g\text{O}_2$  kinetics in the solid phase.

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(37) [sensitizer]  $\sim 5.4 \times 10^{-3}$  M (ref 31).

(38) By using published values for the oxygen diffusion coefficient (D) in PMMA (ref 28 and 29), we calculate a mean oxygen displacement of  $\sim 200$  Å in 100  $\mu\text{s}$  (ref 36). D in a liquid phase solvent is approximately three orders of magnitude larger (ref 27, pp 57–58).

(39) Since the rate constant for  $^1\Delta_g\text{O}_2$  appearance ( $\tau_{1/2} \sim 10$   $\mu\text{s}$ , air saturated sample) must equal the total rate constant for the disappearance of its precursor (ref 19), the excited state sensitizer lives long enough to influence the  $^1\Delta_g\text{O}_2$  decay.

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(32)  $^1\Delta_g\text{O}_2$  sensitizers were irradiated with a pulsed laser at 355 nm (1.5 mJ/pulse) (ref 19, 25, 26, 33).

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(35) The time required for our signal amplitude to change under different ambient oxygen concentrations coincides with the expected rate of oxygen diffusion (ref 38) into the bulk polymer. This result, and other evidence (ref 36), indicates that  $^1\Delta_g\text{O}_2$  phosphorescence originates primarily from the sample interior.

(36) Details will be forthcoming.